Quantum Walks

Kaniuar Bacho

Ruhr University Bochum

July 4, 2023

Kaniuar Bacho	(Ruhr University Bochun	
---------------	-------------------------	--

1/26



- Classical Random Walks
- Quantum Walks
- Application: Grover



Classical Random Walks

3 Quantum Walks

Application: Grover

Kaniuar Bacho	(Ruhr University Bochum)
---------------	--------------------------

3/26

Mathematician George Pólya liked to take morning walks through the woods. He noticed that he would regularly bump into the same couple.

Then he asked himself "what is the probability that these two randomly walking groups bump into each other?"

We need a mathematical framework to answer such questions.

4/26

Mathematician George Pólya liked to take morning walks through the woods. He noticed that he would regularly bump into the same couple.

Then he asked himself "what is the probability that these two randomly walking groups bump into each other?"

We need a mathematical framework to answer such questions.



[https://en.wikipedia.org/wiki/Random_walk] Brownian motion of a dust particle in the room can be modeled by a random walk.

Mathematician George Pólya liked to take morning walks through the woods. He noticed that he would regularly bump into the same couple.

Then he asked himself "what is the probability that these two randomly walking groups bump into each other?"

We need a mathematical framework to answer such questions.



[https://en.wikipedia.org/wiki/Random_walk] Brownian motion of a dust particle in the room can be modeled by a random walk.



[https://de.wikipedia.org/wiki/Random_Walk] Black-Scholes model: stock market prices evolve according to a random walk.

Shuffling a deck of cards can be modeled by a random walk.

Fun fact: 7 *riffle shuffles* are sufficient to consider the deck as "good shuffled". For *overhand shuffle* we need 10000.



[https://betandbeat.com/poker/rules/shuffling/]

5/26

Shuffling a deck of cards can be modeled by a random walk.

Fun fact: 7 *riffle shuffles* are sufficient to consider the deck as "good shuffled". For *overhand shuffle* we need 10000.



[https://betandbeat.com/poker/rules/shuffling/]



Google's search algorithm *PageRank*: measure the importance of a webpage using random surfer model.

Olassical Random Walks

3 Quantum Walks

Application: Grover

	(Ruhr University Bochum)
--	--------------------------

6/26

Classical Random Walks

• **Informally**: a random walk is a series of steps through a state space, where the location of the next step is chosen randomly from among a set of available options.



7/26

Classical Random Walks

• **Informally**: a random walk is a series of steps through a state space, where the location of the next step is chosen randomly from among a set of available options.



- We can define random walks on very different state spaces:
 d dimensional lattice Z^d, graphs, mathematical groups, Riemannian manifolds,...
- We can have discrete-time or continuous-time (even more exotic index sets).
- We can add many features (e.g. absorbing states, reflecting states,...).

1-dimensional simple random walk on $\ensuremath{\mathbb{Z}}$

Initial position $S_0 \in \mathbb{Z}$. After each time step: go to the right with probability p or to the left with 1 - p. Let S_n be the position after n steps. Model it mathematically by:

 $X_i \stackrel{\text{iid}}{\sim} \text{Bernoulli}(p)$

with $Pr[X_i = 1] = p$ and $Pr[X_i = -1] = 1 - p$. Hence our position after *n* steps is

$$S_n = S_{n-1} + X_n$$

= $S_0 + X_1 + X_2 + \dots + X_n$.

• Example:

How does S_n evolve over time for $S_0 = 0$ and p = 1/2?

[YouTube - PBS Infinite Series - What is a Random Walk?]



 S_1 after one time step

Kaniuar Bacho	(Ruhr University Bochum)
---------------	--------------------------

・ロト ・回ト ・ヨト



[YouTube - PBS Infinite Series - What is a Random Walk?]

 S_1 after one time step

-7 -6 -5 -4 -3 -2 -1 0 1 2 3 4 5 6 7 1/4 1/2 1/4 probability probability

 S_2 after two time steps



 S_1 after one time step

two turns 1/4

 S_2 after two time steps



 S_3 after three time steps

9/26



• Now: random walks on finite, *d*-regular, connected graphs G = (V, E) without selfloops.

- Now: random walks on finite, *d*-regular, connected graphs G = (V, E) without selfloops.
- State space is V. Let N := |V|.
 Probability going from vertex x to y is given by

$$\Pr[x, y] = \begin{cases} 1/d, \text{ if } \{x, y\} \in E\\ 0, \text{ else} \end{cases}$$

Encode this information into the $N \times N$ transition matrix P with $P[x, y] = \Pr[x, y]$.

- Now: random walks on finite, *d*-regular, connected graphs G = (V, E) without selfloops.
- State space is V. Let N := |V|.
 Probability going from vertex x to y is given by

$$\Pr[x, y] = \begin{cases} 1/d, \text{ if } \{x, y\} \in E\\ 0, \text{ else} \end{cases}$$

Encode this information into the $N \times N$ transition matrix P with $P[x, y] = \Pr[x, y]$.

Initial distribution: unlike before, starting point X₀ is not determined.
 It is chosen from a probability distribution X₀ ~ p₀, i.e. p₀ is a probability vector over the vertices V.

- Now: random walks on finite, *d*-regular, connected graphs G = (V, E) without selfloops.
- State space is V. Let N := |V|. Probability going from vertex x to y is given by

$$\Pr[x, y] = \begin{cases} 1/d, \text{ if } \{x, y\} \in E\\ 0, \text{ else} \end{cases}$$

Encode this information into the $N \times N$ transition matrix P with $P[x, y] = \Pr[x, y]$.

- Initial distribution: unlike before, starting point X₀ is not determined.
 It is chosen from a probability distribution X₀ ~ p₀, i.e. p₀ is a probability vector over the vertices V.
- We are interest in how X_n evolves over time. Let's compute X_1 :

$$\Pr[X_1 = y] = \sum_{x} \Pr[X_0 = x] \cdot \Pr[x, y] = (p_0 \cdot P)[y].$$

Hence

 $X_1 \sim p_0 \cdot P \Longrightarrow X_n \sim p_0 \cdot P^n.$

• We showed $p_n = p_0 \cdot P^n$. How does it behave in the long term?

- We showed $p_n = p_0 \cdot P^n$. How does it behave in the long term?
- A probability vector π is called *stationary distribution* if $\pi = \pi \cdot P$.

- We showed $p_n = p_0 \cdot P^n$. How does it behave in the long term?
- A probability vector π is called *stationary distribution* if $\pi = \pi \cdot P$.

• In our case above $\pi = (\frac{1}{N}, \frac{1}{N}, ..., \frac{1}{N})$ is a stationary distribution:

$$(\pi \cdot P)[y] = \sum_{x} \pi[x] \cdot \Pr[x, y] = \sum_{x} \frac{1}{N} \cdot \Pr[x, y] = \sum_{x \sim y} \frac{1}{N} \cdot \Pr[x, y] = \frac{1}{N}.$$

- We showed $p_n = p_0 \cdot P^n$. How does it behave in the long term?
- A probability vector π is called *stationary distribution* if $\pi = \pi \cdot P$.
- In our case above $\pi = (\frac{1}{N}, \frac{1}{N}, ..., \frac{1}{N})$ is a stationary distribution:

$$(\pi \cdot P)[y] = \sum_{x} \pi[x] \cdot \Pr[x, y] = \sum_{x} \frac{1}{N} \cdot \Pr[x, y] = \sum_{x \sim y} \frac{1}{N} \cdot \Pr[x, y] = \frac{1}{N}.$$

Uniqueness and Convergence

i) Stationary distribution π is unique (here it is used that G is connected). ii) If G is additionally non-bipartite, then $p_n \xrightarrow{n \to \infty} \pi \quad \forall p_0$.

11/26

 The ε-mixing time measures how fast we approach the stationary distribution with an error of at most ε:

$$\mathrm{MT}(\epsilon) := \min\{n : \sup_{P_0} \|p_0 \cdot P^n - \pi\| \leq \epsilon\}.$$

 The ε-mixing time measures how fast we approach the stationary distribution with an error of at most ε:

$$\mathrm{MT}(\epsilon) := \min\{n : \sup_{p_0} \|p_0 \cdot P^n - \pi\| \le \epsilon\}.$$

• Since *P* is symmetric and real, we have *N* real eigenvalues $\lambda_1 \ge \lambda_2 \ge ... \ge \lambda_N$. Using the Perron–Frobenius theorem we get $1 = \lambda_1 > \lambda_2 \ge ... \ge \lambda_N > -1$.

 The ε-mixing time measures how fast we approach the stationary distribution with an error of at most ε:

$$\mathrm{MT}(\epsilon) := \min\{n : \sup_{P_0} \|p_0 \cdot P^n - \pi\| \le \epsilon\}.$$

- Since *P* is symmetric and real, we have *N* real eigenvalues $\lambda_1 \ge \lambda_2 \ge ... \ge \lambda_N$. Using the Perron–Frobenius theorem we get $1 = \lambda_1 > \lambda_2 \ge ... \ge \lambda_N > -1$.
- Define the spectral gap $\delta := 1 \max_{i \ge 2} |\lambda_i| > 0$.

 The ε-mixing time measures how fast we approach the stationary distribution with an error of at most ε:

$$\mathrm{MT}(\epsilon) := \min\{n : \sup_{P_0} \|p_0 \cdot P^n - \pi\| \le \epsilon\}.$$

- Since P is symmetric and real, we have N real eigenvalues λ₁ ≥ λ₂ ≥ ... ≥ λ_N. Using the Perron–Frobenius theorem we get 1 = λ₁ > λ₂ ≥ ... ≥ λ_N > −1.
- Define the spectral gap $\delta := 1 \max_{i \ge 2} |\lambda_i| > 0$.

Convergence rate

$$\operatorname{MT}(\epsilon) \leq \frac{1}{\delta} \cdot (\frac{1}{2}\log(N) + \log\left(\frac{1}{\epsilon}\right))$$

12/26

Let M ⊆ V be a subset of the vertices, which are called *marked*.
 Assume we start with the stationary distribution and want to hit an element in M.
 The *hitting time* asks for how fast we can find such a marked element

$$\begin{aligned} \tau_{\mathcal{M}} &:= \min_{t \geq 0} \{ t \mid X_t \in \mathcal{M} \land X_0 \sim \pi \}, \\ & \text{HT}_{\mathcal{M}} := \mathbb{E}[\tau_{\mathcal{M}}]. \end{aligned}$$

Let M ⊆ V be a subset of the vertices, which are called *marked*.
 Assume we start with the stationary distribution and want to hit an element in M.
 The *hitting time* asks for how fast we can find such a marked element

$$\tau_{\mathcal{M}} := \min_{t \ge 0} \{ t \mid X_t \in \mathcal{M} \land X_0 \sim \pi \},$$
$$\operatorname{HT}_{\mathcal{M}} := \mathbb{E}[\tau_{\mathcal{M}}].$$

Complexity bound

$$\mathrm{HT}_{M} \in O\left(\frac{1}{\delta} \cdot \frac{N}{|M|}\right)$$

Kaniuar Bacho	(Ruhr University Bochum)
---------------	--------------------------

Let M ⊆ V be a subset of the vertices, which are called *marked*.
 Assume we start with the stationary distribution and want to hit an element in M.
 The *hitting time* asks for how fast we can find such a marked element

$$\tau_{\mathcal{M}} := \min_{t \ge 0} \{ t \mid X_t \in \mathcal{M} \land X_0 \sim \pi \},$$
$$\operatorname{HT}_{\mathcal{M}} := \mathbb{E}[\tau_{\mathcal{M}}].$$

Complexity bound

$$\mathrm{HT}_{M} \in O\left(\frac{1}{\delta} \cdot \frac{N}{|M|}\right)$$

• This random walk search algorithm can find marked elements in $O\left(\frac{1}{\delta} \cdot \frac{N}{|M|}\right)$. Start at vertex x, then repeat until find marked vertex: check if current vertex is marked if not, run a random walk for roughly $1/\delta$ steps to get close to the uniform distribution.

Can we improve this by switching to a quantum setting?

13/26

Classical Random Walks

Quantum Walks

Application: Grover

Kaniuar Bacho	(Ruhr University Bochum)
---------------	--------------------------

- How to convert random walk into a quantum walk?
- Naive approach would be a unitary U of the form $|x\rangle \mapsto \frac{1}{\sqrt{d}} \sum_{y \sim x} |y\rangle$. But this cannot exist in general: consider a graph with two different vertices x and
 - z, which have the same neighbours. So $U\ket{x}=U\ket{z}$, hence $\ket{x}=\ket{z}$. ${}_{z}$

- How to convert random walk into a quantum walk?
- Naive approach would be a unitary U of the form $|x\rangle \mapsto \frac{1}{\sqrt{d}} \sum_{y \sim x} |y\rangle$. But this cannot exist in general: consider a graph with two different vertices x and z, which have the same neighbours. So $U|x\rangle = U|z\rangle$, hence $|x\rangle = |z\rangle \cdot \frac{4}{2}$
- Classical random walk on Z: coin flip followed by a shift. **Solution**: combine them into one object. Add *coin space* to the *state space* and define two unitaries.

Think of it as: position of the walker & direction that the walker points to.

Quantum Walks: quantizing random walks (part 2)

• **On graphs**: write $|x, y\rangle = |x\rangle |y\rangle$ for being at vertex x pointing to y. This can be considered to quantum walk over the edges (x, y)!

- On graphs: write $|x, y\rangle = |x\rangle |y\rangle$ for being at vertex x pointing to y. This can be considered to quantum walk over the edges (x, y)!
- Shift: if the walker moves to y, what should be the new direction? As we came from x, it should be x. This is implemented by the SWAP-operation $S |x, y\rangle = |y, x\rangle$.

- On graphs: write $|x, y\rangle = |x\rangle |y\rangle$ for being at vertex x pointing to y. This can be considered to quantum walk over the edges (x, y)!
- Shift: if the walker moves to y, what should be the new direction? As we came from x, it should be x. This is implemented by the SWAP-operation S |x, y > |y, x >.

• Coin toss: set
$$|\psi_x\rangle = \frac{1}{\sqrt{d}} \sum_{y \sim x} |x, y\rangle$$
. Define the coin toss by

$$C = 2\left(\sum_{x \in V} |\psi_x\rangle \langle \psi_x|\right) - I,$$

which is a reflection around $|\psi_x\rangle$.

- On graphs: write $|x, y\rangle = |x\rangle |y\rangle$ for being at vertex x pointing to y. This can be considered to quantum walk over the edges (x, y)!
- Shift: if the walker moves to y, what should be the new direction? As we came from x, it should be x. This is implemented by the SWAP-operation S |x, y > |y, x >.

• Coin toss: set
$$|\psi_x\rangle = \frac{1}{\sqrt{d}} \sum_{y \sim x} |x, y\rangle$$
. Define the coin toss by

$$C = 2\left(\sum_{x \in V} |\psi_x\rangle \langle \psi_x|\right) - I,$$

which is a reflection around $|\psi_x\rangle$.

• Let the quantum walk operator be $W = S \cdot C$.

16/26

Let *H_P* be the Hilbert space spanned by the positions of the walker. On the line it is spanned by {|*i*⟩ : *i* ∈ ℤ}. Add the coin space *H_C* to it, spanned by {|0⟩, |1⟩} = {|→⟩, |←⟩}. States of the total system are in the space *H* = *H_C* ⊗ *H_P*.

- Let *H_P* be the Hilbert space spanned by the positions of the walker. On the line it is spanned by {|*i*⟩ : *i* ∈ ℤ}. Add the coin space *H_C* to it, spanned by {|0⟩, |1⟩} = {|→⟩, |←⟩}. States of the total system are in the space *H* = *H_C* ⊗ *H_P*.
- The conditional translation is described by

$$\begin{split} S &= \left| \rightarrow \right\rangle \left\langle \rightarrow \right| \otimes \sum_{i} \left| i + 1 \right\rangle \left\langle i \right| \\ &+ \left| \leftarrow \right\rangle \left\langle \leftarrow \right| \otimes \sum_{i} \left| i - 1 \right\rangle \left\langle i \right| \end{split}$$

So $| \rightarrow \rangle \otimes | i \rangle$ is transformed to $| \rightarrow \rangle \otimes | i + 1 \rangle$ and $| \leftarrow \rangle \otimes | i \rangle$ to $| \leftarrow \rangle \otimes | i - 1 \rangle$.

- Let *H_P* be the Hilbert space spanned by the positions of the walker. On the line it is spanned by {|*i*⟩ : *i* ∈ ℤ}. Add the coin space *H_C* to it, spanned by {|0⟩, |1⟩} = {|→⟩, |←⟩}. States of the total system are in the space *H* = *H_C* ⊗ *H_P*.
- The conditional translation is described by

$$\begin{split} S &= \left| \rightarrow \right\rangle \left\langle \rightarrow \right| \otimes \sum_{i} \left| i + 1 \right\rangle \left\langle i \right| \\ &+ \left| \leftarrow \right\rangle \left\langle \leftarrow \right| \otimes \sum_{i} \left| i - 1 \right\rangle \left\langle i \right| \end{split}$$

So $| \rightarrow \rangle \otimes | i \rangle$ is transformed to $| \rightarrow \rangle \otimes | i + 1 \rangle$ and $| \leftarrow \rangle \otimes | i \rangle$ to $| \leftarrow \rangle \otimes | i - 1 \rangle$.

• As a coin flip, we will take the Hadamard coin $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$.

- Let *H_P* be the Hilbert space spanned by the positions of the walker. On the line it is spanned by {|*i*⟩ : *i* ∈ ℤ}. Add the coin space *H_C* to it, spanned by {|0⟩, |1⟩} = {|→⟩, |←⟩}. States of the total system are in the space *H* = *H_C* ⊗ *H_P*.
- The conditional translation is described by

$$egin{aligned} \mathcal{S} &= \left|
ightarrow
ight
angle \left<
ightarrow \left|
ightarrow \sum_{i} \left| i + 1
ight
angle \left< i
ight| \ &+ \left|
ightarrow
ight
angle \left<
ightarrow \left|
ightarrow \sum_{i} \left| i - 1
ight
angle \left< i
ight| \end{aligned}$$

So $| \rightarrow \rangle \otimes | i \rangle$ is transformed to $| \rightarrow \rangle \otimes | i + 1 \rangle$ and $| \leftarrow \rangle \otimes | i \rangle$ to $| \leftarrow \rangle \otimes | i - 1 \rangle$.

- As a coin flip, we will take the Hadamard coin $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$.
- Start with state $|\Psi_0\rangle = |\rightarrow\rangle \otimes |0\rangle$, apply $H \otimes I$ to get $\frac{1}{\sqrt{2}} |\rightarrow\rangle \otimes |0\rangle + \frac{1}{\sqrt{2}} |\leftarrow\rangle \otimes |0\rangle$.
 - Apply S to get $|\Psi_1\rangle = \frac{1}{\sqrt{2}} | \rightarrow \rangle \otimes | 1 \rangle + \frac{1}{\sqrt{2}} | \leftarrow \rangle \otimes | -1 \rangle.$

• Measure to get $\{ | \rightarrow \rangle \otimes | 1 \rangle, | \leftarrow \rangle \otimes | -1 \rangle \}$ each with probability 1/2. If we continue and measure after each iteration, we obtain the plain classical random walk on \mathbb{Z} .

- Measure to get $\{|\rightarrow\rangle \otimes |1\rangle, |\leftarrow\rangle \otimes |-1\rangle\}$ each with probability 1/2. If we continue and measure after each iteration, we obtain the plain classical random walk on \mathbb{Z} .
- What does happen if we do not measure in-between? It will behave completely different!

- Measure to get $\{|\rightarrow\rangle \otimes |1\rangle, |\leftarrow\rangle \otimes |-1\rangle\}$ each with probability 1/2. If we continue and measure after each iteration, we obtain the plain classical random walk on \mathbb{Z} .
- What does happen if we do not measure in-between? It will behave completely different!
- Let $U = S \cdot (H \otimes I)$, then we get

$$egin{aligned} |\Psi_3
angle &= U^3 \left|\Psi_0
ight
angle \ &= rac{1}{2\sqrt{2}} \left(|{
ightarrow}
angle \otimes |3
angle + |{
ightarrow}
angle \otimes |1
angle + 2 \left|{
ightarrow}
angle \otimes |1
angle - |{
ightarrow}
angle \otimes |-1
angle + |{
ightarrow}
angle \otimes |-3
angle
ight). \end{aligned}$$

So $|3\rangle$, $|-3\rangle$, $|-1\rangle$ appear with probability 1/8, but $|1\rangle$ with 5/8.

18/26

Classical random walk vs. quantum walk after 50 steps:



[https://en.wikipedia.org/wiki/Quantum_walk]

Reason for drifting to the right: the Hadamard coin treats $|\rightarrow\rangle$ and $|\leftarrow\rangle$ differently. $|\leftarrow\rangle$ will appear more often with negative and positive coefficients, which will cancel out.

• Remember
$$|\psi_x\rangle = \frac{1}{\sqrt{d}} \sum_{y \sim x} |x, y\rangle$$
, $C = 2\left(\sum_{x \in V} |\psi_x\rangle \langle \psi_x|\right) - I$ and $W = S \cdot C$.

æ

イロン イ団 とく ヨン イヨン

• Remember
$$|\psi_x\rangle = \frac{1}{\sqrt{d}} \sum_{y \sim x} |x, y\rangle$$
, $C = 2\left(\sum_{x \in V} |\psi_x\rangle \langle \psi_x|\right) - I$ and $W = S \cdot C$.

• Our *quantum walk algorithm* will work similar to Grover's algorithm. Define two states

$$|G\rangle := rac{1}{\sqrt{|M|}} \sum_{x \in M} |\psi_x\rangle, \ |B\rangle := rac{1}{\sqrt{N - |M|}} \sum_{x \notin M} |\psi_x\rangle.$$

And

$$|\pi\rangle := |U\rangle = \frac{1}{\sqrt{N}} \sum_{x} |\psi_x\rangle.$$

• Remember
$$|\psi_x\rangle = \frac{1}{\sqrt{d}} \sum_{y \sim x} |x, y\rangle$$
, $C = 2\left(\sum_{x \in V} |\psi_x\rangle \langle \psi_x|\right) - I$ and $W = S \cdot C$.

• Our *quantum walk algorithm* will work similar to Grover's algorithm. Define two states

$$|G\rangle := rac{1}{\sqrt{|M|}} \sum_{x \in M} |\psi_x\rangle, \ |B\rangle := rac{1}{\sqrt{N - |M|}} \sum_{x \notin M} |\psi_x\rangle.$$

And

$$|\pi\rangle := |U\rangle = rac{1}{\sqrt{N}} \sum_{x} |\psi_{x}\rangle \,.$$

• Now we only have to implement the reflection through $|B\rangle$ and through $|U\rangle,$ like in Grover.

The first one is easy: just check if the first register has a marked element. If yes, put a -1.

The second one is harder. Here we will use the quantum walk W.

20/26

• Remember
$$|\psi_x\rangle = \frac{1}{\sqrt{d}} \sum_{y \sim x} |x, y\rangle$$
, $C = 2\left(\sum_{x \in V} |\psi_x\rangle \langle \psi_x|\right) - I$ and $W = S \cdot C$.

• Our *quantum walk algorithm* will work similar to Grover's algorithm. Define two states

$$|G\rangle := rac{1}{\sqrt{|M|}} \sum_{x \in M} |\psi_x\rangle, \ |B
angle := rac{1}{\sqrt{N - |M|}} \sum_{x \notin M} |\psi_x
angle.$$

And

$$|\pi
angle := |U
angle = rac{1}{\sqrt{N}} \sum_{x} |\psi_{x}
angle \,.$$

• Now we only have to implement the reflection through $|B\rangle$ and through $|U\rangle,$ like in Grover.

The first one is easy: just check if the first register has a marked element. If yes, put a -1.

The second one is harder. Here we will use the quantum walk W.

• We have that
$$|\pi\rangle = \frac{1}{\sqrt{Nd}} \sum_{(x,y)\in E} |x,y\rangle$$
 is a stationary state of W .

Szegedy's spectral lemma (2004)

The non-trivial spectrum of the quantum walk operator W can be characterized as follows: for every eigenvalue $\lambda_j = \cos \theta_j \neq 1$ of P, the quantum walk operator W has a pair of eigenvalues $e^{\pm i\theta_j}$ with $\theta_j \in (0, \pi)$.

We have a correspondence between the eigenvalues of *P* and of *W* by projecting them to the circle. The *phase gap* is defined by $\Delta := \min\{\theta_j \mid \theta_j \neq 0\}$. By the lemma we have $\theta_j = \arccos(\lambda_j)$. Hence $\Delta = \arccos(\lambda_2)$.



Szegedy's spectral lemma (2004)

The non-trivial spectrum of the quantum walk operator W can be characterized as follows: for every eigenvalue $\lambda_i = \cos \theta_i \neq 1$ of P, the quantum walk operator W has a pair of eigenvalues $e^{\pm i\theta_j}$ with $\theta_i \in (0, \pi)$.



• For the analysis(!) use spectral decomposition for any state $|\Psi\rangle = \alpha_0 |\pi\rangle + \sum \alpha_j |v_j\rangle$ with $W |v_j\rangle = e^{i\theta_j} |v_j\rangle$ and $\theta_j \ge \Delta > 0$.

・ロト ・四ト ・ヨト ・ヨト

- For the analysis(!) use spectral decomposition for any state $|\Psi\rangle = \alpha_0 |\pi\rangle + \sum \alpha_j |v_j\rangle$ with $W |v_j\rangle = e^{i\theta_j} |v_j\rangle$ and $\theta_j \ge \Delta > 0$.
- Use quantum phase estimation QPE to distinguish between these two summands. The first one has eigenvalue with angle 0, the other one with $\theta_i \ge \Delta > 0$.

- For the analysis(!) use spectral decomposition for any state $|\Psi\rangle = \alpha_0 |\pi\rangle + \sum \alpha_j |v_j\rangle$ with $W |v_j\rangle = e^{i\theta_j} |v_j\rangle$ and $\theta_j \ge \Delta > 0$.
- Use quantum phase estimation QPE to distinguish between these two summands. The first one has eigenvalue with angle 0, the other one with $\theta_j \ge \Delta > 0$.
- Use QPE with Δ precision. Store the angle in a new register. Now a apply the unitary $2 \cdot I \otimes |0\rangle \langle 0| I$ to add sign -1 if angle is not 0.

- For the analysis(!) use spectral decomposition for any state $|\Psi\rangle = \alpha_0 |\pi\rangle + \sum \alpha_j |v_j\rangle$ with $W |v_j\rangle = e^{i\theta_j} |v_j\rangle$ and $\theta_j \ge \Delta > 0$.
- Use quantum phase estimation QPE to distinguish between these two summands. The first one has eigenvalue with angle 0, the other one with $\theta_i \ge \Delta > 0$.
- Use QPE with Δ precision. Store the angle in a new register. Now a apply the unitary 2 · *I* ⊗ |0⟩ ⟨0| − *I* to add sign -1 if angle is not 0.
- Uncompute the QPE again by applying the inverse (make sure we have no entanglements). We arrive at $|\Psi'\rangle = \alpha_0 |\pi\rangle \sum \alpha_j |v_j\rangle$.

- For the analysis(!) use spectral decomposition for any state $|\Psi\rangle = \alpha_0 |\pi\rangle + \sum \alpha_j |v_j\rangle$ with $W |v_j\rangle = e^{i\theta_j} |v_j\rangle$ and $\theta_j \ge \Delta > 0$.
- Use quantum phase estimation QPE to distinguish between these two summands. The first one has eigenvalue with angle 0, the other one with $\theta_i \ge \Delta > 0$.
- Use QPE with Δ precision. Store the angle in a new register. Now a apply the unitary $2 \cdot I \otimes |0\rangle \langle 0| I$ to add sign -1 if angle is not 0.
- Uncompute the QPE again by applying the inverse (make sure we have no entanglements). We arrive at $|\Psi'\rangle = \alpha_0 |\pi\rangle \sum \alpha_j |v_j\rangle$.
- We implemented $2 \cdot |\pi\rangle \langle \pi| I$ successfully. This is our reflection around $|\pi\rangle$.

- For the analysis(!) use spectral decomposition for any state $|\Psi\rangle = \alpha_0 |\pi\rangle + \sum \alpha_j |v_j\rangle$ with $W |v_j\rangle = e^{i\theta_j} |v_j\rangle$ and $\theta_j \ge \Delta > 0$.
- Use quantum phase estimation QPE to distinguish between these two summands. The first one has eigenvalue with angle 0, the other one with $\theta_i \ge \Delta > 0$.
- Use QPE with Δ precision. Store the angle in a new register. Now a apply the unitary $2 \cdot I \otimes |0\rangle \langle 0| I$ to add sign -1 if angle is not 0.
- Uncompute the QPE again by applying the inverse (make sure we have no entanglements). We arrive at $|\Psi'\rangle = \alpha_0 |\pi\rangle \sum \alpha_j |v_j\rangle$.
- We implemented $2 \cdot |\pi\rangle \langle \pi| I$ successfully. This is our reflection around $|\pi\rangle$.
- All in all our quantum walk searching algorithm needs

$$O\left(rac{1}{\Delta}\cdot\sqrt{rac{N}{|M|}}
ight)=O\left(\sqrt{rac{1}{\delta}\cdotrac{N}{|M|}}
ight)$$

Classical Random Walks

3 Quantum Walks

Application: Grover

イロン イロン イヨン イヨン

Convert problem into graph

- High level idea how to solve problems using quantum walks:
 - 1) Convert problem P into a graph G with marked elements M, s.t.
 - "Finding marked element in G" \Leftrightarrow "Finding solution of P".
 - 2) Now quantum walk on G until we hit M.

Convert problem into graph

- High level idea how to solve problems using quantum walks:
 1) Convert problem P into a graph G with marked elements M, s.t.
 "Finding marked element in G" ⇔ "Finding solution of P".
 2) Now quantum walk on G until we hit M.
- Solve *unstructured search problem* using quantum walks (instead of Grover). Given an *N*-bit string with exactly *T* 1's. We want to find a 1. Convert to *N*-complete graph with *N* vertices. It should be complete because we have the freedom to query any index in the list, meaning we can jump to any other position/vertex.

Now the marked elements corresponds to the 1's, meaning |M| = T. Moreover one can show $\delta = 1 - \frac{1}{N-1} \approx 1$. Hence the complexity to find a marked element is

$$O\left(\sqrt{\frac{N}{|M|}}\right) = O\left(\sqrt{\frac{N}{T}}\right).$$

References

- [1] Bhattacharya, Rabindra Nath, and Edward C. Waymire. *Random walk, Brownian motion, and martingales.* Springer, 2021.
- [2] Levin, David A., and Yuval Peres. *Markov chains and mixing times*. Vol. 107. American Mathematical Soc., 2017.
- [3] Doyle, Peter G., and J. Laurie Snell. *Random walks and electric networks*. Vol. 22. American Mathematical Soc., 1984.
- [4] Lovász, László. "Random walks on graphs." Combinatorics, Paul erdos is eighty 2.1-46 (1993): 4.
- [5] De Wolf, Ronald. "Quantum computing: Lecture notes." arXiv preprint arXiv:1907.09415 (2019).
- [6] Wong, Thomas G. "Unstructured search by random and quantum walk." arXiv preprint arXiv:2011.14533 (2020).
- [7] Kempe, Julia. "Quantum random walks: an introductory overview." Contemporary Physics 44.4 (2003): 307-327.
- [8] Apers, Simon. "Lecture 1: Random walks and quantum walks" (2022).

Thank you for your attention :-) Questions?

・ロト ・四ト ・ヨト ・ヨト